

EUROPEAN CENTRE FOR THEORETICAL STUDIE IN NUCLEAR PHYSICS AND RELATED AREAS

TRENTO, ITALY al Member of the ESF Expert Committee NuPECC

Information and Statistics in Nuclear Experiment and Theory ISNET-3

Trento, November 16-20, 2015

Density Functional Theory meets Bayesian Neural Networks: A New Paradigm in the Study of Neutron Stars



Jorge Piekarewicz - Florida State University

Highlights of 2015



Welcome to JPhysG's 2015 highlights! Before you get to the articles, here are a few of my personal choices.



First, the way the community reacted and engaged with our focus issue, Enhancing the interaction between nuclear experiment and theory through information and statistics was outstanding. Linking theory with experiment is vital for any field and I look forward to seeing more research on the topic in both nuclear and particle physics.



Bayesian Methods in Nuclear Physics INT Program 2016

2017 ICNT Program: Extracting Bulk Properties of Neutron-Rich Matter with Transport Models in Bayesian Perspective. March 22 — April 12, 2017 at FRIB/MSU

Neutron Stars: Very Few Historical Facts

- Chandrasekhar shows that massive stars will collapse (1931)
- Chadwick discovers the neutron (1932) (... predicted earlier by Majorana but never published)
- Baade-Zwicky introduce the concept of a neutron star (1933)

(... Landau mentions dense stars that look like giant nuclei!)

- Solution Oppenheimer-Volkoff use GR to compute the structure of neutron stars (1939) (... predict $M_{\star} \simeq 0.7 M_{\odot}$ as maximum neutron star mass)
- Jocelyn Bell discovers neutron stars (1967)









Neutron Stars: Unique Cosmic Laboratories

- Neutron stars are the remnants of massive stellar explosions (CCSN)
 - Bound by gravity NOT by the strong force
 - Catalyst for the formation of exotic state of matter
 - Satisfy the Tolman-Oppenheimer-Volkoff equation (v_{esc} /c ~ 1/2)
- Only Physics that the TOV equation is sensitive to: Equation of State
 EOS must span about 11 orders of magnitude in baryon density
- Increase from 0.7→ 2 M_{sun} transfers ownership to Nuclear Physics!
 Predictions on stellar radii differ by several kilometers!



$$\frac{dM}{dr} = 4\pi r^2 \mathcal{E}(r)$$

$$\frac{dP}{dr} = -G \frac{\mathcal{E}(r)M(r)}{r^2} \left[1 + \frac{P(r)}{\mathcal{E}(r)}\right]$$

$$\left[1 + \frac{4\pi r^3 P(r)}{M(r)}\right] \left[1 - \frac{2GM(r)}{r}\right]^{-1}$$

Need an EOS: $P = P(\mathcal{E})$ relation Nuclear Physics Critical



The Composition of the Outer Crust Enormous sensitivity to nuclear masses

- System unstable to cluster formation
 - BCC lattice of neutron-rich nuclei imbedded in e-gas
- Composition emerges from relatively simple dynamics
- Competition between electronic and symmetry energy

$$E/A_{\rm tot} = M(N,Z)/A + \frac{3}{4}Y_e^{4/3}k_{\rm F} + \text{lattice}$$



Both - for neutron-star crusts and r-process nucleosynthesis







INTERNATIONAL JOURNAL OF HIGH-ENERGY PHYSICS



DFT meets BNN

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Use DFT to predict nuclear masses
Train BNN by focusing on residuals *M(N,Z) = M_{DFT}(N,Z) + \delta M_{BNN}(N,Z)*

- Systematic scattering greatly reduced
- Predictions supplemented by theoretical errors



The Equation of State of Neutron-Rich Matter

- The EOS of asymmetric matter: $\alpha = (N-Z)/A$; $x = (\rho \rho_0)/3\rho_0$; T = 0
 - $\rho_0 \simeq 0.15 \text{ fm-3} \text{saturation density} \leftrightarrow \text{nuclear density}$ $\mathcal{E}(\rho, \alpha) \simeq \mathcal{E}_0(\rho) + \alpha^2 \mathcal{S}(\rho) \simeq \left(\epsilon_0 + \frac{1}{2}K_0 x^2\right) + \left(J + Lx + \frac{1}{2}K_{\text{sym}} x^2\right) \alpha^2$
- Symmetric nuclear matter saturates:
 - $\epsilon_0 \simeq -16 \text{ MeV} \text{binding energy per nucleon} \leftrightarrow \text{nuclear masses}$
 - $K_0 \simeq 230 \text{ MeV} \text{nuclear incompressibility} \leftrightarrow \text{nuclear "breathing" mode}$
- Density dependence of symmetry poorly constrained:
 - Solution Symmetry energy ↔ masses of neutron-rich nuclei
 - Subscript{S







Model Building: The Protocol



PHYSICAL REVIEW C 90, 044305 (2014) Building relativistic mean field models for finite nuclei and neutron stars

Wei-Chia Chen^{*} and J. Piekarewicz[†] Department of Physics, Florida State University, Tallahassee, Florida 32306, USA



$$\mathcal{L}_{\text{Yukawa}} = \bar{\psi} \left[g_{\text{s}} \phi - \left(g_{\text{v}} V_{\mu} + \frac{g_{\rho}}{2} \tau \cdot \mathbf{b}_{\mu} + \frac{e}{2} (1 + \tau_3) A_{\mu} \right) \gamma^{\mu} \right] \psi$$
$$\mathcal{L}_{\text{self}} = \frac{\kappa}{3!} (g_{\text{s}} \phi)^3 - \frac{\lambda}{4!} (g_{\text{s}} \phi)^4 + \frac{\zeta}{4!} g_{\text{v}}^4 (V_{\mu} V^{\mu})^2 + \Lambda_{\text{v}} \left(g_{\rho}^2 \, \mathbf{b}_{\mu} \cdot \mathbf{b}^{\mu} \right) \left(g_{\text{v}}^2 V_{\nu} V^{\nu} \right)$$

Nuclear Density Functional Theory (DFT)

- Ab-initio calculations of heavy nuclei remains daunting task
- Search for energy functional valid over a large physics domain

"from finite nuclei to neutron stars"

- Incorporate physics insights into the construction of the functional
- Accurately calibrated to various properties of finite nuclei masses, charge radii, and giant monopole resonances
- Empirical constants encode physics beyond mean field
- Sempirical constants obtained from the optimization of a quality measure

Nucleus	Observable	Experiment	NL3	FSU	FSU2
¹⁶ O	B/A	7.98	8.06	7.98	8.00
	$R_{ m ch}$	2.70	2.75	2.71	2.73
⁴⁰ Ca	B/A	8.55	8.56	8.54	8.54
	$R_{ m ch}$	3.48	3.49	3.45	3.47
⁴⁸ Ca	B/A	8.67	8.66	8.58	8.63
	$R_{ m ch}$	3.48	3.49	3.48	3.47
⁶⁸ Ni	B/A	8.68	8.71	8.66	8.69
	$R_{ m ch}$	—	3.88	3.88	3.86
⁹⁰ Zr	B/A	8.71	8.70	8.68	8.69
	$R_{ m ch}$	4.27	4.28	4.27	4.26
¹⁰⁰ Sn	B/A	8.25	8.30	8.24	8.28
	$R_{\rm ch}$	_	4.48	4.48	4.47
¹¹⁶ Sn	B/A	8.52	8.50	8.50	8.49
	$R_{\rm ch}$	4.63	4.63	4.63	4.61
¹³² Sn	B/A	8.36	8.38	8.34	8.36
	$R_{\rm ch}$	4.71	4.72	4.74	4.71
¹⁴⁴ Sm	B/A	8.30	8.32	8.32	8.31
	$R_{\rm ch}$	4.95	4.96	4.96	4.94
²⁰⁸ Pb	B/A	7.87	7.90	7.89	7.88
	R_{ch}	5.50	5.53	5.54	5.51

Nucleus	TAMU	RCNP	NL3	FSU	FSU2
⁹⁰ Zr	17.81 ± 0.35	_	18.76	17.86	17.93 ± 0.09
¹¹⁶ Sn	15.90 ± 0.07	15.70 ± 0.10	17.19	16.39	16.47 ± 0.08
144 Sm	15.25 ± 0.11	15.77 ± 0.17	16.29	15.55	15.59 ± 0.09
²⁰⁸ Pb	14.18 ± 0.11	13.50 ± 0.10	14.32	13.72	13.76 ± 0.08

Bayes' Theorem: Application to Model Building



- QCD is the fundamental theory of the strong interactions!
 - M: A theoretical MODEL with parameters and biases
 - D: A collection of experimental and observational DATA
- The Prior P(M): An insightful transformation in DFT $(g_{s}, g_{v}, g_{\rho}, \kappa, \lambda, \Lambda_{v}) \iff (\rho_{0}, \epsilon_{0}, M^{*}, K, J, L)$

Solution The Likelihood
$$P(D|M) = \exp(-\chi^2/2)$$

 $\chi^2(D,M) = \sum_{n=1}^N \frac{\left(O_n^{(\text{th})}(M) - O_n^{(\exp)}(D)\right)^2}{\Delta O_n^2}$

The Marginal Likelihood; overall normalization factor

Heaven and Earth The enormous reach of the neutron skin

- Neutron-star radii are sensitive to the EOS near $2\rho_0$
- Neutron star masses sensitive to EOS at much higher density
- Neutron skin correlated to a host of neutron-star properties
 - Stellar radii, proton fraction, enhanced cooling, moment of inertia
- Neutron skin of heavy nuclei and NS radii driven by same physics
 Difference in length scales of 18 orders of magnitude!!







Searching for L: The Strategy

- Establish a powerful physical argument connecting L to R_{skin}
 - Where do the extra 44 neutrons in ²⁰⁸Pb go? Competition between surface tension and the *difference* $S(\rho_0)$ - $S(\rho_{surf}) \simeq L$. *The larger the value of L, the thicker the neutron skin of* ²⁰⁸Pb
- Ensure that "your" DFT supports the correlation
- Ensure that "all" accurately-calibrated DFT support the correlation (... "all models are equal but some models are more equal than others")



Electroweak Measurement of Neutron Densities

- PREX@JLAB: First electroweak (clean!) evidence in favor of Rskin in Pb
- Precision hindered by radiation issues
 - Excellent control of systematic uncertainties
 - Statistical uncertainties 3 times larger than promised
- PREX-II and CREX to run in 2018
- Original goal of 1% in neutron radius

$$A_{\rm PV} = \frac{G_F Q^2}{2\sqrt{2}\pi\alpha} \left[\underbrace{1 - 4\sin^2\theta_W}_{\approx 0} - \frac{F_n(Q^2)}{F_p(Q^2)} \right]$$

- Neutral weak-vector boson Z_0 couples preferentially to neutrons
- PV provides a clean measurement of neutron densities (and R_n)

	up-quark	down-quark	proton	neutron			
γ -coupling	+2/3	-1/3	+1	0			
Z_0 -coupling	pprox +1/3	pprox -2/3	pprox 0	-1			
$g_{ m v}\!=\!2t_{z}-4Q\sin^{2} heta_{ m W}\!pprox\!2t_{z}\!-\!Q$							







The incompressibility of neutron rich matter: Why is tin so fluffy?

Workshop on Nuclear Incompressibility

University of Notre Dame July 14-15, 2005

The Joint Institute for Nuclear Astrophysics (JINA) will organize a 2-day Workshop focused on Nuclear Incompressibility and the Nuclear Equation of State, to be held at the University of Notre Dame during July 14-15, 2005.

This meeting follows a similar Workshop held at Notre Dame in January 2001, and the Symposium on Nuclear Equation of State used in Astrophysics Models, held at the ACS meeting in Philadelphia last Summer.

The primary aim of the Workshop is to bring together interested physicists from the areas of Astrophysics, Giant Resonances, and Heavy-Ion Reactions, to discuss current status of experiments and theoretical models related to nuclear incompressibility and the equation of state, and to explore what experiments might be needed to clarify some of the outstanding issues.

Most of the Workshop will be devoted to talks, with a lot of time allowed for discussions and interactions. In that spirit, we will follow a somewhat flexible schedule for the talks.

There is no registration fee but participants are requested to register via the webpage (www.jinaweb.org), so that we can make appropriate arrangements.

For further information, please contact: Kathy Burgess (kburgess@nd.edu)

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R.

JINA



The Joint Institute for Nuclear Astrophysics May 18, 2005

Outcome: A window into L through systematic measurements of the GMR across a long isotopic chain



Available online at www.sciencedirect.com ScienceDirect



Nuclear Physics A 788 (2007) 36c-43c

The Giant Monopole Resonance in the Sn Isotopes: Why is Tin so "Fluffy"?

U. Garg,^a T. Li,^a S. Okumura,^b H. Akimune^c M. Fujiwara,^b M.N. Harakeh,^d H. Hashimoto,^b M. Itoh,^e Y. Iwao,^f T. Kawabata,^g K. Kawase,^b Y. Liu,^a R. Marks,^a T. Murakami,^f K. Nakanishi,^b B.K. Navak,^a P.V. Madhusudhana Rao,^a H. Sakaguchi,^f Y. Terashima,^f M. Uchida,^h Y. Yasuda,^f M. Yosoi,^b and J. Zenihiro^f



Isotopic Dependence of the Giant Monopole Resonance in the Even-A ^{112–124}Sn Isotopes and the Asymmetry Term in Nuclear Incompressibility

T. Li,¹ U. Garg,¹ Y. Liu,¹ R. Marks,¹ B. K. Nayak,¹ P. V. Madhusudhana Rao,¹ M. Fujiwara,² H. Hashimoto,² K. Kawase,² K. Nakanishi,² S. Okumura,² M. Yosoi,² M. Itoh,³ M. Ichikawa,³ R. Matsuo,³ T. Terazono,³ M. Uchida,⁴ T. Kawabata,⁵ H. Akimune,⁶ Y. Iwao,⁷ T. Murakami,⁷ H. Sakaguchi,⁷ S. Terashima,⁷ Y. Yasuda,⁷ J. Zenihiro,⁷ and M. N. Harakeh⁸

PHYSICAL REVIEW C 86, 024303 (2012)

Giant monopole resonances and nuclear incompressibilities studied for the zero-range and separable pairing interactions

P. Veselý,^{1,*} J. Toivanen,¹ B. G. Carlsson,² J. Dobaczewski,^{1,3} N. Michel,¹ and A. Pastore⁴

Onwards and upwards to GMRs in unstable nuclei!

 $K_0(\alpha) = K_0 + K_\tau \alpha^2;$ $K_\tau = K_{\text{sym}} - \underbrace{6L} + \dots$

Acceived 10 May 2007, Publishe

Nicroscopic linear response.

Description of the Siant mont

PHYSICAL REVIEW C 75, 031301(R) (2007)

* Florida 32305. USA

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1011 COUDING

With is the equation of state for tin so soft?

PHYSICAL REVIEW C 18 US4304 (3008)

PHYSICAL REVIEW C 39, 0343406 (2)

Electric Dipole Polarizability







IVGDR: The quintessential nuclear excitation

- Out-of-phase oscillation of neutrons vs protons Symmetry energy acts as restoring force
- Energy weighted sum rule largely model independent
- Inverse energy weighted sum strongly correlated to L
 Actually ... Jα_D strongly correlated to L
 Important contribution from Pygmy resonance
- High quality data emerging from RCNP, GSI, HIGS On a variety of nuclei such as Pb, Sn, Ni, Ca, ... and hopefully in the future along isotopic chains



Neutron-Star Radii

PRL 115, 161101 (2015) P

week ending 16 OCTOBER 2015

Compactness of Neutron Stars

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- Guillot et al., assume all neutron stars share a common radius Assumption in MR observable rather than on the EOS
- One-to-one correspondence between EOS and MR TOV equation + EOS — Unique MR relation
- Lindblom's inversion algorithm shows the inverse also true! [APJ 398, 569 (1992)] TOV equation + MR — Unique Equation of State
- For a given "common" radius MR profile examine whether: Resulting EOS is causal or superluminal for stellar masses below $2M_{\odot}$
- For a given "common" radius MR profile, to prevent causality violations Stellar radius of a $1.4M_{\odot}$ must exceed 10.7 km!





"We have detected gravitational waves; we did it" David Reitze, February 11, 2016



The dawn of gravitational wave astronomy

- Initial black hole masses are 36 and 29 solar masses
- Final black hole mass is 62 solar masses;

3 solar masses radiated in GW!





What will we learn from neutron-star mergers







Tidal polarizability scales as R⁵!







NS radius measured to better than 1km!

What *else* will we learn from neutron-star mergers

How were the heavy elements from iron to uranium made?

- LIGO will provide critical insights into the behavior of ultra dense matter
- Merger rate and ejecta mass unknown Galactic merger rate depends on EOS: 4x10⁻⁵ (soft) 4x10⁻⁴ (stiff) per year to account for observation



Soft: Rn-Rp is small → neutron star more compact merger is more violent → higher abundance





The Origin of the Solar System Elements

My Collaborators

My FSU Collaborators

- Genaro Toledo-Sanchez
- Karim Hasnaoui
- Bonnie Todd-Rutel
- Brad Futch
- Jutri Taruna
- Farrukh Fattoyev
- Wei-Chia Chen
- Raditya Utama



My Outside Collaborators

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- W. Nazarewicz (MSU)
- N. Paar (U. Zagreb)
- M.A. Pérez-Garcia (U. Salamanca)
- P.G.- Reinhard (U. Erlangen-Nürnberg)
- X. Roca-Maza (U. Milano)
- D. Vretenar (U. Zagreb)





